

Electrical Measurements of Organic Semiconductors

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Overview



Introducing Faro Organic Semiconductors Electrical Measurements DLTS Interface states Field Effect Transistors

Introducing Faro



Introducing Faro

Geography





Universidade do Algarve Faro, Portugal





OptoEl



http://www.ualg.pt/fct/adeec/optoel

OPTO-ELECTRONICS





CEOT



Centro de Electrónica, Opto-electrónica e Telecomunicações



Center for Electronics, Opto-electronics and Telecommunications

Funding





Projects



Projects in which OptoEl is involved:

- SELOA (European TMR project): organic electronics
- BioFET: measuring living cells with organic based FETs
- TNT: designing sensors for detecting landmines
- MONA-LISA (European RTD project): single-molecule organic FETs
- bilateral project with Bologna: electronics based on BDT

MONA-LISA





MONA-LISA (<u>Mo</u>lecular <u>Nanostructures made by</u> <u>Lithography and their</u> <u>Scaling Aspects</u>)

Study of organic materials in nano-scale devices.

Ultimately: single molecule electronics

Why? Organics have low mobility due to disorder. Single molecule has high order, hence high µ.

Organic Semiconductors



Organic Semiconductors

Organic Semiconductors



They are cheap to make No high purity required (no clean rooms, etc) No extreme conditions required (*T*, *p*)
Taylor-made: Infinite number of possibilities bandgap engineering
Optical efficiency high in complete color range
Flexible. Imagine the roll-up screen you are looking at now. It could be the display of the future.

Disadvantages:

- Most unstable in air (so far!)
- Slow (low μ)! Forget application of fast electronics.

Applications: LEDs, displays (background ill.), solar cells, tagging ...

Examples







• Every semiconducting polymer has a "backbone" of undercoordinated carbon atoms

• 4^{th} electron is in weak $p_z p_z$ bonds. Loosely bound -> metal • Jahn-Teller deformation of backbone: creation of alternating single and double bonds Material Band gap

$-\mathbf{CH} = \mathbf{CH} - \mathbf{CH} = \mathbf{CH} -$

- This causes opening of a band gap
 -> semiconductor
- band gap $\pm 2.5 \text{ eV}$
- wide band gap ¹/₂con

Material	Band gap
SiO ₂	>10 eV
C (diamond)	5.47 eV
GaN	3.36 eV
Polymers	2.5 eV
GaAs	1.42 eV
Si	1.12 eV
Ge	0.66 eV

3D conduction



Conduction along backbone is easy

Conduction from chain to chain is Difficult. Hopping Conduction

Two activation energies for conduction



The **stacking of the molecules** is very important! Best 3D conductivity achieved by vacuum deposition

Electrical measurements



Electrical measurements

Note: a survey of electrical measurement techniques is available at our internet pages: http://www.ualg.pt/fct/adeec/optoel/theory

Electrical Measurements



<u>Devices</u>:

- bulk
- Schottky barrier
- pn-junction
- FETS



"Lucky for us there's an electrical outlet right here!"

<u>Techniques</u>:

- current-voltage (DC)
- capacitance, conductance (AC)
- admittance spectroscopy
- Transient techniques:
 - capacitance transients
 - DLTS
 - TSC

Information:

- conduction model
- carrier type
- carrier mobility
- shallow levels

 position
 density

 deep levels

 position
 density

 dielectric constant
 barrier height

Each of these techniques are standard characterization techniques for classical semiconductors (Si, GaAs, GaN, etc), but care has to be taken to apply them to organic materials!

Schottky barrier





- metal and ¹/₂ con have different Fermi level
- electrons will flow from metal to ¹/₂con
- build-up of (space) charge Q (uncompensated ionized acceptors)
- causes electric field and voltage drop (band bending, $V_{\rm bi}$)
- over a range W (depletion width)

$$V_{\rm bi} = \chi + V_{\rm n} - \phi_{\rm m}$$

Schottky barrier: Capacitance





slope reveals N_A
extrapolation reveals V_{bi}

Schottky barrier: Admittance spectroscopy





Schottky barrier: Admittance spectroscopy





Maximum at $1/\omega_{\text{max}} = R_b C_b(C_b + C_d)$

$$R_{\rm b} \sim \exp(-E_{\rm a}/kT)$$

We can determine the bulk activation energy from the tan δ data





example



Schottky barrier: Admittance spectroscopy







Measuring above ω_{max} will give no information about the interface C_d , R_d !

Most commercial equipment works at 1 MHz. Most organic materials cut around 10 kHz.

BDT; oFET



BDT, the "light emitting FET"



For a light-emitting FET we need to be able to inject both electrons and holes; the device has to be ambipolar. BDT has shown to have this behavior.









DLTS: "Change the bias and let's look how the capacitance evolves over time"

Capacitance depends on bias (remember, something like C~1/V^{1/2})
A new depletion width has to be reached. At the end:

 $C = \varepsilon A/W$

- For shallow levels: response is immediate. Limited only by speed at which free carriers can move out (μ_p) .
- For deep levels: the charges have to come off there first.





- 1. Free holes move out of interface region. Immediate increase of W (C \downarrow)
- 2. This creates a region where the deep levels are off-equilibrium
- 3. Charges are slowly emitted from the deep levels there
 - higher space-charge density
 - less depletion width is needed to reach condition $ff_0^W \rho(x) d^2x = V_{bb}$
- 4. W slowly shrinks again a little. Increased capacitance







Monitoring τ over temperature will give us E_a

Very sensitive and very accurate!





 $1/\tau = e_{\rm p} = \gamma T^2 \sigma \exp(-E_{\rm a}/kT)$

MEH-PPV on Silicon2 minority traps: a, c1 majority trap: b



First successful "DLTS" experiment in an organic semiconductor P. Stallinga *et al.* J. Appl. Phys. **89**, 1713 (2001)



Standard DLTS procedure:

- Of the entire set of data, take only two points, at t_1 and t_2
- The DLTS signal is then $S = C(t_1) C(t_2)$

For low-T: $\tau = 0$, $C(t_1) = C(t_2)$: S = 0For high-T: $\tau = 0$, $C(t_1) = C(t_2)$: S = 0

Maximum when: $\tau_0 T^{-2} \exp(E_a/kT) = (t_2 - t_1)/[\ln(t_1/t_2)]$





... not followed (don't throw away data)



DLTS is

- very easy to perform. "Walk-away" measurements
- sensitive
- reliable data with accurate energy determination
- "Fingerprint" spectra of defects
- able to determine density of defects. $\Delta C/C = N_T/2(N_A N_D)$

Problems for organic semiconductors

- cut-off frequency (10 kHz)
- deep levels slow transients (T=300 K, $\tau=1$ min)
- condition $N_T \ll N_A$ not true in most cases (transients non exponential)

Field Effect Transistors



Field Effect Transistors

Field Effect Transistors





FET: sexi-thiophene (α T6)





Mobility changes as function of gate voltage

$\alpha T6$ transfer curve





Perfect straight line when plotted as 6th-root of current vs. Vg

 $I_{ds} \sim (V_g - V_t)^6$

$\alpha T6$ transfer curve





Plotted as 5th-root. Consistent with saturation-region data

 $I_{ds} \sim V_{ds} (V_g - V_t)^5$

Mobility





Gate-voltage-dependent mobility up to 1x10⁻³ cm²/Vs

Variable Range Hopping



PHYSICAL REVIEW B

VOLUME 57, NUMBER 20

15 MAY 1998-II

Theory of the field-effect mobility in amorphous organic transistors

M. C. J. M. Vissenberg* Philips Research Laboratories, 5656 AA Eindhoven, The Netherlands and Instituut–Lorentz, University of Leiden, 2300 RA Leiden, The Netherlands

M. Matters[†] Philips Research Laboratories, 5656 AA Eindhoven, The Netherlands (Received 13 January 1998)

The field-effect mobility in an organic thin-film transistor is studied theoretically. From a percolation model of hopping between localized states and a transistor model an analytic expression for the field-effect mobility is obtained. The theory is applied to describe the experiments by Brown *et al.* [Synth. Met. **88**, 37 (1997)] on solution-processed amorphous organic transistors, made from a polymer (polythienylene vinylene) and from a small molecule (pentacene). Good agreement is obtained, with respect to both the gate voltage and the temperature dependence of the mobility. [S0163-1829(98)01320-4]

Gate-voltage dependent mobility predicted and measured before

JOURNAL OF APPLIED PHYSICS

VOLUME 85, NUMBER 6

15 MARCH 1999

Gate voltage dependent mobility of oligothiophene field-effect transistors

Gilles Horowitz,⁴⁰ Riadh Hajlaoui, Denis Fichou, and Ahmed El Kassmi Laboratoire des Matériaux Moléculaires, CNRS, 2 rue Henry-Dunant, 94320 Thiais, France

(Received 5 May 1998; accepted for publication 8 December 1998)

Organic field-effect transistors, in which the active semiconductor is made of oligothiophenes of various lengths, have been fabricated and characterized. A method is developed to estimate the field-effect mobility μ corrected for the contact series resistance. The mobility is found to increase by a factor of nearly 100 from quaterthiophene (4T) to octithiophene (8T). More importantly, μ increases quasilinearly with gate voltage. The origin of this gate bias dependence is discussed. One explanation could be the presence of traps that limit charge transport. Alternatively, the gate-voltage dependence is tentatively attributed to a dependence of the mobility with the concentration of carriers in the accumulation layer. \bigcirc 1999 American Institute of Physics. [S0021-8979(99)02506-2]

Variable-range hopping

Variable Range Hopping





VRH vs. experiment $\alpha T6$





VRH vs. experiment





Title





Each dot represents the value of the slope in a transfer curve at a different temperature (of previous page)

Slopes ' α ' not well predicted by VRH

Threshold voltage





Impurity levels freeze out in F10 and not in X20

T6 - FET mobility conclusions



Bad morphology Failure of VRH theory V_T shows huge number of traps

> Better morphology Hints of VRH theory V_T shows smaller number of traps

> > Good morphology Standard FET theory V_T shows nearly no traps

Other FET measurements





Scanning speed



Hysteresis



Pulsed / CW





Interface states

Interface states







Interface states studied in Schottky barriers of

Aluminum - Terrylene

Interface states are important because they can cause (un)wanted voltage drops at the interface and thus block/facilitate carrier injection .





When the Fermi level is resonant with the interface states we see a contribution to the capacitance and inductance.

But only when the AC frequency is low enough

With the bias we select the bunch of interface states measured









The depth of the levels we can find from the temperature dependence.





The number of interface states we can find from the height of the peak in conductance.



1: select the interface states with the bias

2: measure their activation energy via the temperature dependence 3: measure the density of states at E_A from the spectrum





Experiment









Experiment







Experiment



 $\omega_{\rm max} = 1/\tau = \exp\left(-E_a/kT\right)$





First direct electrical observation of interface states in organic semiconductors, P. Stallinga *et al.* Organic Electronics, 2001.

Summary



Electrical Measurements in Faro:

DLTS in Si/MEH-PPV

Interface states in Terrylene/Al

Field Effect Transistors: transport model

10 Q for your attention